# The Neutron and X-ray Crystal Structure of (Na,Li), Th<sub>6</sub>F<sub>31</sub> with More Than One Least-Squares Minimum\*

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The compound  $(Na, Li)_7 Th_6 F_{31}$  crystallizes in space group P3c1 with  $a_0 = 9.9056 \pm 0.0003$  and  $c_0 = 13.2820 \pm 0.0005$  Å at 25 °C. The calculated density is 6.045 g.cm<sup>-3</sup> and Z=2. There are 18 atoms in the asymmetric unit. Two independent Th<sup>4+</sup> ions and F(1) through F(10) are on the general position 6(d); Th(1) is coordinated by 9F<sup>-</sup> and Th(2) is coordinated by 10F<sup>-</sup>. Ion F(11) is on position 2(a). The three sodium ions Na(1), Na(2) and Na(3) occupy positions 2(c), 2(b) and 2(a) respectively. The ion Na(3) is surrounded by 12F<sup>-</sup> and Na(2) and Na(1) are surrounded by 6F<sup>-</sup> at distances less than 3.0 Å. Sodium (1) and Na(2) have five and three more F<sup>-</sup> nearest neighbors respectively between 3.0 and 3.5 Å. Lithium(1) is octahedrally coordinated at position 2(a). Lithium(2) is at general position 6(d) and is surrounded by 7F<sup>-</sup> nearest neighbors less than 3.0 Å distant.

There are at least two least-squares minima for each of two sets of data obtained with X-rays and neutrons. The ranges of cation-anion interatomic distances for the originally determined X-ray and neutron minima are respectively: Th(1)-F,  $2\cdot23-2\cdot62$ ,  $2\cdot28-2\cdot52$  Å; Th(2)-F,  $2\cdot25-2\cdot63$ ,  $2\cdot29-2\cdot57$  Å; Na(1)-F,  $2\cdot26-3\cdot54$ ,  $2\cdot26-3\cdot41$  Å; Na(2)-F,  $2\cdot25-3\cdot34$ ,  $2\cdot30-3\cdot07$  Å; Na(3)-F,  $2\cdot56-2\cdot96$ ,  $2\cdot50-3\cdot00$  Å; Li(1)-F,  $2\cdot03-2\cdot06$ ,  $1\cdot96-2\cdot06$  Å; Li(2)-F,  $1\cdot77-2\cdot89$ ,  $1\cdot92-2\cdot77$  Å. Refinement of the X-ray parameters with neutron data, and *vice versa*, did not change the parameters significantly; two very significant minima remained, with differences as large as  $30\sigma$ .

### Introduction

The ternary system NaF-LiF-ThF<sub>4</sub> and its intermediate compounds are being investigated as part of the Molten Salt Breeder Reactor Program at the Oak Ridge National Laboratory (Thoma, Insley, Landau, Friedman & Grimes, 1959; Thoma, Insley, Hebert, Friedman & Weaver, 1963). The ternary intermediate compound (Na,Li)<sub>7</sub>Th<sub>6</sub>F<sub>31</sub> was regarded as the terminus of a solid solution join with the high temperature modification of NaThF<sub>5</sub> ( $\alpha$ -NaThF<sub>5</sub>, Brunton, Insley, McVay & Thoma, 1965) and was erroneously reported as NaLi<sub>2</sub>Th<sub>2</sub>F<sub>11</sub>. Both phases crystallize with the same space group and have similar unit-cell dimensions.

### Experimental

We decided to gather both neutron and X-ray diffraction data for the crystal structure determination. The positions of the thorium atoms and most of the fluorine atoms could be determined from the three-dimensional Patterson sections computed from both sets of data. The best positional parameters and temperature factors for the low atomic number atoms would come from the neutron data.

Single crystals of  $(Na,^7Li)_7Th_6F_{31}$  (hereafter referred to as NLT) were selected from an ingot of the composition: NaF-20 mole %, <sup>7</sup>LiF-40 mole %, ThF<sub>5</sub>-40

mole % which was grown in a modified Stockbarger-Bridgman gradient furnace (Weaver, Ross, Sturm, Eorgan & Thoma, 1964). The crystals were tumbleground in an abrasive-lined ball mill until they were roughly cubic with rounded edges. These crystals were etched smooth in warm concentrated aqueous  $Al(NO_3)_3$ solution and were thermally shocked by repeated immersions alternately in liquid nitrogen and mineral oil at 20 °C. The crystals were translucent after this treatment. One of us (DRS) observed a greenish-blue thermochromism in the crystals at 77 °K which may be the result of a phase transition at low temperature. Fragments of the same material were submitted for chemical analysis, the results of which are reported in Table 1.

Thirty-three high angle X-ray reflections were used to refine the unit-cell parameters by least-squares analysis;  $a_0 = 9.9056 \pm 0.003$  and  $c_0 = 13.2820 \pm 0.0005$  Å at 25 °C. ( $\lambda$ , Cu  $K\alpha_1 = 1.54050$ , Cu  $K\alpha_2 = 1.54434$  Å). The space group is P3c1 and the conditions limiting possible reflections are:  $h\bar{h}0l$ , l=2n present only. Piezoelectric tests on the Mellon Institute Giebe–Scheibe detector were negative but the final structure, from both X-ray and neutron data, is noncentrosymmetrical.

The calculated density is 6.045 g.cm<sup>-3</sup>, the measured density is 5.94 g.cm<sup>-3</sup>, Z=2 and the linear absorption,  $\mu = 1735.51$  cm<sup>-1</sup> for Cu K $\alpha$  radiation, and is less than 0.1 for 1.0770 Å neutrons.

### Neutron determination

The crystal used for neutron measurements was about 2mm on a side, and weighed 10.4 mg, somewhat small

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Table 1. Results of chemical analyses of (Na, 7Li)<sub>7</sub>Th<sub>6</sub>F<sub>31</sub> crystals

	7Li	Na	Th	F
Mean %	1.16	4.38	66.72	28.22
Standard deviation for individual	0.11	0.55	0.26	0.31
Number of atoms $Th = 12$	6.92	7.95	12	61.99

Table 2. Tostional parameters and temperature factors for (1(a,L))711613	Table 2.	Positional	parameters	and	temperature	factors	for	(Na,Li)7Th6F31
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ATOM		x	Y	z	$\beta_{11}^{a}$	β22	β33	β1 2	β13	β23
Th(1)	NO <sup>h</sup> XN XO NX	0.0782(4) <sup>b</sup> 0.0786(2) 0.0847(2) 0.0850(4)	0.3954(4) 0.3954(2) 0.4094(2) 0.4088(4)	0.0 <sup>°</sup> 0.0 0.0 0.0	0.0012(3) 0.0035(2) 0.0044(2) 0.0016(4)	0.0012(3) 0.0029(2) 0.0039(3) 0.0013(4)	0.0006(2) 0.0005(1) 0.0007(1) 0.0003(2)	0.0006(3) 0.0016(2) 0.0024(2) 0.0010(3)	0.0002(2) 0.0000(1) 0.0000(2) -0.0004(2)	-0.0001(2) 0.0000(1) -0.0003(1) -0.0002(2)
Th(2)	NO XN XO NX	0.9148(4) 0.9153(2) 0.9214(2) 0.9220(4)	0.5911(4) 0.5906(2) 0.6045(2) 0.6043(4)	0.2059(2) 0.2060(2) 0.2061(2) 0.2059(2)	0.0008(3) 0.0033(3) 0.0036(2) 0.0011(4)	0.0017(3) 0.0035(2) 0.0035(2) 0.0013(4)	0.0003(2) 0.0006(1) 0.00035(9) 0.0006(2)	0.0007(3) 0.0020(1) 0.0017(2) 0.0005(3)	0.0001(2) 0.0002(1) 0.0000(1) 0.0002(2)	0.0003(2) 0.0000(2) 0.0002(1) -0.0001(2)
F(1)	NO XN XO	0.1370(7) 0.144(4) 0.139(4)	0.4291(8) 0.443(4) 0.441(4)	-0.1862(5) -0.192(3) -0.181(3)	0.0029(7) 0.005(2) 0.002(1)	0.0038(7) d d	0.0007(3)	0.0003(6)	-0.0005(4)	-0.0001(4)
		0.1335(4)	0.4435(9)	-0.1803(6)	0.0036(8)	0.0041(8)	0.0008(3)	0.0010(7)	0.0002(4)	0.0006(4)
F(2)	XN XO NX	0.8658(8) 0.869(3) 0.864(4) 0.8615(8)	0.5552(7) 0.564(3) 0.563(4) 0.5692(8)	0.3853(6) 0.376(3) 0.388(3) 0.3916(6)	0.0036(7) 0.003(2) 0.002(1) 0.0025(7)	0.0030(7) d d 0.0037(8)	0.0009(3)	0.0007(6)	-0.0008(4)	-0.0002(4)
	NO	0.1191(8)	0 3061(8)	-0.3759(6)	0.0023(7)	0.0030(8)	0.0007(5)	0.0001(0)	-0.0002(4)	0.0003(4)
F(3)	XN XO NX	0.114(4) 0.129(5) 0.1197(9)	0.306(3) 0.315(4) 0.3088(9)	-0.374(3) -0.374(3) -0.3770(6)	0.005(2) 0.003(2) 0.0028(7)	d d 0.0041(9)	0.0016(4)	0.0018(7)	0.0004(5)	-0.0009(5)
	NO	0.8781	0.6893(8)	0.5827(6)	0.0037(8)	0.0028(7)	0.0012(4)	0.0015(6)	0.0006(4)	-0.0004(4)
P(4)	XN XO NX	0.874(4) 0.886(5) 0.8784(9)	0.682(4) 0.690(4) 0.6912(9)	0.586(3) 0.587(3) 0.5811(6)	0.005(2) 0.003(2) 0.0028(8)	d d 0,0027(8)	0.0015(4)	0.0013(6)	0.0012(5)	0.0003(5)
	NO	0.3431(6)	0.5103(7)	-0.0444(5)	0.0015(6)	0.0047(7)	0.0022(3)	0,0001(5)	-0.0005(3)	0.0005(4)
F(5)	XN XO NX	0.340(4) 0.357(5) 0.3512(8)	0.517(4) 0.459(5) 0.4488(8)	-0.053(3) -0.062(3) -0.0555(6)	0.006(2) 0.004(2) 0.0031(7)	d d 0.007(1)	0.0025(4)	0.0004(7)	-0.0003(4)	-0.0006(5)
	NO	0.6479(8)	0.5494(7)	0.2615(5)	0.0069(9)	0.0038(7)	0.0021(3)	0.0006(7)	0.0004(4)	0.0003(4)
F(6)	XO NX	0.668(4) 0.6564(7)	0.541(4) 0.493(5) 0.4906(8)	0.258(3) 0.249(3) 0.2501(6)	0.008(2) 0.004(2) 0.0016(6)	d d 0.0054(8)	0.0021(3)	-0.0001(6)	-0.0007(4)	0.0007(4)
	NO	0.1033(8)	0.1966(7)	-0.0655(6)	0.0059(8)	0.0025(7)	0,0015(3)	0.0031(7)	-0.0001(4)	-0.0006(4)
<b>P</b> (7)	XO NX	0.079(4) 0.0808(8)	0.196(4) 0.1893(8)	-0.061(3) -0.0588(7)	0.003(2) 0.0009(9) 0.0033(7)	d 0.0010(7)	0.0024(4)	0.0000(6)	-0.0001(4)	-0.0010(4)
F(8)	NO XN XO	0.9191(7) 0.918(3) 0.893(4)	0.8111(7) 0.812(3) 0.803(3)	0.2643(7) 0.257(2) 0.267(3)	0.0014(6) 0.002(2) 0.0003(9)	0.0031(7) d d	0.0018(3)	0.0005(6)	0.0009(4)	0.0001(4)
	NX	0.8985(9)	0.8058(8)	0.2712(6)	0.0053(8)	0.0038(8)	0,001?(3)	0.0039(7)	-0.0006(4)	-0.0010(5)
₽(9)	NO XN XO	0.1295(8) 0.125(4) 0.117(4)	0.5851(7) 0.584(4) 0.576(4)	-0.3531(6) -0.364(3) -0.347(3)	0.0026(7) 0.007(2) 0.001(1)	0.0020(7) d d	0.0017(4)	0.0003(6)	-0.0014(4)	0.0004(4)
	NA	0.117(1)	0.5787(8)	-0.3459(6)	0.0049(9)	0.0014(7)	0.0014(3)	0.0007(6)	-0.0012(5)	0.0004(4)
<b>P</b> (10)	XN XO NX	0.8836(8) 0.879(3) 0.868(5) 0.8716(9)	0.4214(8) 0.421(3) 0.417(4) 0.4153(8)	0.5508(6) 0.543(3) 0.555(4) 0.5578(6)	0.0022(7) 0.004(2) 0.002(1) 0.0029(8)	0.0033(7) d d 0.0035(8)	0.0013(3)	0.0003(6)	-0.0007(4)	0.0008(4)
	NO	0.66667	0.33333	=0.306(1)	0.0038(7)	•	0.008(1)	*	-0.0015(5)	•
F(11)	XN XO	0.66667	0.33333	-0.321(5) -0.321(8)	0.008(3)	å		-	-	-
	NX	0.66667	0.33333	-0.309(2)	0.017(3)	Ĩ	0.011(3)	f	f	r
Na(l)	NO XN XO	0.66667 0.66667 0.66667	0.33333 0.33333 0.33333	-0.047(2) -0.045(4) -0.055(4)	0.002(1) 0.003(2) 0.001(2)	f d d	0.003(2)	f	£	f
	NX	0.66667	0.33333	-0.054(2)	0.005(2)	f	0.0005(8)	f	f	f
Na (2)	NO XN XO NX	0.33333 0.33333 0.33333 0.33333	0.66667 0.66667 0.66667	0.258(2) 0.254(4) 0.248(4) 0.251(2)	0.005(2) 0.005(2) 0.001(1)	f d d	0.0006(8)	1	t	ſ
	NO	0.55555	0.00007	0.231(2)	0.0003(9)		0.003(1)	1	1	1
Na(3)	XN XO NX	0.0	0.0	0.121(2) 0.118(3) 0.084(5) 0.083(2)	0.006(2) 0.005(2) 0.009(3) 0.006(2)	r d d	0.002(2)	T T	r •	r •
	NO	0.0	0.0	-0.139(3)	0.003(2)	-	0.002(2)	-	-	•
L1(1)	XN XO NX	0.0 0.0 0.0	0.0 0.0 0.0	-0.19(2) -0.15(2) -0.15(3)	0.02(1) 0.008(6) 0.004(2)	d d f	0.000(2)	r t	r f	r f
	NO	0.233(4)	0.283(3)	-0.188(3)	0,027(6)8	0.004(3)	0.014(4)	-0.011(4)	0.012(4)	-0.001(3)
L1(2)	XN XO NX	0.28(2) 0.272(9)* 0.286(8)	0.29(2) 0.310(9) 0.250(7)	-0.20(2) -0.221(8) -0.13(2)	0.01(1) 0.003(3) 0.04(2) <sup>B</sup>	d d 0.019(7)	0.07(2)	-0.030(9)	0.04(2)	-0.04(1)

<sup>8</sup>Coefficients in the Temperature Factor: EXP- $(\beta_{11}h^2 + \beta_{22}k^2 + B_{33}k^2 + 2\beta_{13}hk + 2\beta_{23}kk + 2\beta_{12}hk)$ .

<sup>b</sup>Standard Error (in parentheses) Corresponds to Last Significant Digit in Parameter.

CDefines Origin Along Z.

dX-Ray Temperature Factors for F, Na, and Li are Constrained to be Isotropic.

 $f_{Only}$  those  $\beta_{ij}$  Adjusted as Independent Variables are Shown. Symmetry Constraints Upon the  $\beta_{ij}$  of all Atoms in Special Positions are:  $\beta_{11} - \beta_{22} = 2\beta_{12}$  and  $\beta_{13} - \beta_{23} = 0$ . (Levy (1956)).

<sup>g</sup>Temperature factor is not positive-definite.

<sup>b</sup>XO - Original X-ray Parameters and Data.
 NO - Original Neutron Parameters and Data.
 XN - X-ray Data, NO Parameters
 NX - Neutron Data, XO Parameters.

for the purpose. Petrographic examination, neutron and X-ray Laue photographs, and X-ray precession photographs confirmed that the crystal remained unfractured after liquid nitrogen immersion.

Approximately 2200 neutron reflections were measured by the step-scanning method, with the Oak Ridge Automatic Three Circle Neutron Diffractometer at the Oak Ridge Research Reactor (Busing, Smith, Peterson & Levy, 1964) and reduced to a set of intensity data and their corresponding standard errors, by the methods and CDC 1604 Å computer programs described by Brown & Levy (1964). The neutron wavelength was 1.0770 Å. Each measurement comprised 51 steps over a 5° 2 $\theta$  range; each step corresponded to 20,000 monitor counts. A standard reflection (080) was measured after each 50 reflections. We used an unpublished program by C.K.Johnson to form a set of 1394 independent (non-equivalent) intensities, and their standard errors. In this program, the 400 reflection of a standard NaCl crystal served to establish the absolute scale. When combining the standard errors of redundant equivalent reflections, the root mean-square deviation was selected if it was larger than the square root of the sum of the individual variances. A 3% standard error was combined with the statistical standard error to give the total standard error for least-squares weighting.

Positions of two thorium atoms, one lithium atom, and several fluorine atoms were located in a three-dimensional Patterson calculation. The lithium-7 atom, by virtue of its negative coherent scattering length  $(-0.21 \times 10^{-12} \text{ cm})$  was readily observed, in the presence of stronger scatterers, within 0.01 Å of its final refined position. The scattering length for lithium-7 was determined by Peterson & Smith (1967); for the other atoms, we used values tabulated by Bacon (1962): 0.35, 1.01, and  $0.55 \times 10^{-12}$  cm for sodium, thorium, and fluorine, respectively. The scattering lengths were not varied in the least-squares refinement.

#### X-ray determination

The neutron data crystal was crushed and the fragments were ground in an air-driven race. A sphere 130  $\mu$ in diameter was mounted on a computer controlled Picker four-circle single-crystal orienter equipped with a scintillation-counter detector. All of the independent reflections,  $h \ge 0$ ,  $k \ge 0$ ,  $l \ge 0$ , out to  $2\theta = 146^{\circ}$  were measured with Ni-filtered Cu K $\alpha$  radiation by the  $2\theta$  scan technique. Additional equivalent reflections, *hkl*, were measured at high  $2\theta$  angles. The 998 reflection data were corrected for Lorentz-polarization factors and for absorption and were reduced to structure factors,  $F_o$ . The scattering factors for Na+, Li+, Th4+, and F- were taken from Cromer & Waber (1965). The scattering factor for Th<sup>4+</sup> includes the real and imaginary parts of the anomalous dispersion,  $\Delta f' = -4$  and  $\Delta f'' = -15$ electrons (International Tables for X-ray Crystallography 1962). The anomalous dispersion parts of the scattering factors are used by a modified version of the Busing, Martin & Levy (1962) least-squares program. The value of -15 electrons for  $\Delta f''$  is used because the absolute configuration of the crystal was determined by Hamilton's (1965) significance test on the *R* factor to be such that  $\Delta f'' = -15$  electrons gives the best *R* factor (R=0.0600 for  $\Delta f''=15$  electrons and R=0.0596 for  $\Delta f''=-15$  electrons) for a one-dimensional hypothesis and 760 degrees of freedom at the 0.005 significance level. (See Ueki, Zalkin & Templeton (1966) and Cruickshank & McDonald (1967) for a discussion of errors caused by neglect of anomalous scattering.)

A three-dimensional Patterson synthesis gave the positions of the Th<sup>4+</sup> ions at general positions 6(d); Th(1), x=0.09, y=0.41, z=0.0 and Th(2), x=0.91, y=0.59, z=0.20 for starting parameters. The positions of most of the other ions were derived from three-dimensional difference syntheses and by inferences from stereoscopic views of models of the structure.

The atomic positions and temperature factors were adjusted by the method of least-squares with XFLS, a modification of the Busing, Martin & Levy (1962) program. The quantity  $\Sigma w(F_o - F_c)^2$  was minimized and the summation was over all of the reflections. The refinement was based upon  $|F_o|$  for the X-ray data and the corresponding weighting factor was the reciprocal of  $\sigma^2(F_o)$ , where  $\sigma(F_o) = \sigma(F_o^2)/2F_o$ . The empirical equation.

$$\sigma^2(F_0^2) = [T + B + 0.0225(T - B)^2]/[A(Lp)^2],$$

was used to estimate the variance in this case [Lp= Lorentz-polarization, A= absorption correction, T=total and B= background counts (Busing & Levy, 1957)].

## **Discussion of results**

The initial structure models were based upon the assumption that, regardless of the chemical analysis (Table 1) the stoichiometry was probably close to Na<sub>6</sub>Li<sub>6</sub>Th<sub>12</sub>F<sub>60</sub> because of the preliminary results of phase equilibria studies (Thoma, 1966), which placed a compound of that stoichiometry on a join with NaThF<sub>5</sub>. We also assumed that the space group was  $P\overline{3}c1$  because both piezoelectric and pyroelectric tests were negative. Both assumptions proved to be false.

The usual statistical tests of Wilson (1949) and of Howells, Phillips & Rogers (1949) were ambiguous. Pseudo-symmetric pairs F(7,8) and F(9,10) had large isotropic temperature factors with both X-ray and neutron data and the least-square refinement would not converge. When space group P3c1 was tested, the thorium ions remained nearly centrosymmetric but there were large departures from centrosymmetry by most of the fluorine ions, particularly F(7), F(8), F(9)and F(10); F(11) has no pseudo-centrosymmetrical equivalent.

The refinement of the atomic parameters was carried out independently by each of the authors to obtain the final results labeled XO and NO in Table 2. D. P. Shoemaker suggested that the large disparity between these sets of parameters might be the result of the convergence of the least-square refinements to two different minima.

The results of refining the XO parameters with neutron data and NO parameters with X-ray data are listed in Table 2. These parameters XN and NX respectively confirm that there are at least two minima for each of the two sets of data. Additional least-squares refinements were made on the parameters with the Xray data by choosing new starting parameters for the Th ions several standard errors away from XO and NO. The least-squares results either converged on the previous minima or diverged and blew up depending upon how far the new parameters were from the original set. If starting parameters are chosen midway between XO and NO, the least-squares refinement does not converge.

The differences between parameters which refine to the two minima (for each of the two data sets) may be as large as  $30\sigma$  but with the exception of the thermal parameters, F(1) y parameter and Li(2) x and y parameters, the differences between parameters for each set of data within each minimum are  $3\sigma$  or less. The differences between the thermal parameters,  $\beta_{11}$  and  $\beta_{22}$ , for the two independent Th ions obtained from refinement to the same minimum of the neutron and X-ray data are as large as the values obtained in the refinement to the two different minima (Table 2).

The best structure seems to be that represented by the parameters NO (Table 2) from the original neutron data. The discrepancy index R for NO is significantly less than that for NX at the 0.005 significance level and some of the interatomic distances from the XO parameters (Table 2) appear to be unreasonable, *e.g.* Th(1)-F(2) 2.18 Å; F(1)-F(11) 1.72 Å; Li(2)-F(8) 1.77 Å. The NO structure is, therefore, discussed in the rest of the paper.

The final atomic parameters and temperature factors are listed in Table 2. Anisotropic temperature factors were calculated for all of the atoms from the neutron data and for thorium only from X-ray data. The X-ray temperature factors for all of the fluorine, lithium, and sodium ions are isotropic.

We adjusted the occupancy of the Li and Na sites by least squares and they converged to values which did not differ significantly from full, unmixed occupancy.



Fig. 1. The cation contents of one asymetrical unit of  $(Na, Li)_7 Th_6 F_{31}$  showing the coordination polyhedron around each cation. One fourth of a unit cell is outlined.



Fig.2. Coordination polyhedron around Th(1).

The stoichiometry which best fits the crystallographic data is Na<sub>6</sub>Li<sub>8</sub>Th<sub>12</sub>F<sub>62</sub>, contrary to the preliminary report (Sears & Brunton, 1967). This result also is in acceptable agreement with the chemical analysis (Table 1). We prefer to use the less specific formula  $(Na,Li)_7 Th_6 F_{31}$ because there is a group of compounds known to have this stoichiometry: Na<sub>7</sub>Zr<sub>6</sub>F<sub>31</sub> (Burns, Ellison & Levy, 1968); metastable  $Na_7Th_6F_{31}$  and many stable compounds (Thoma, 1962). Furthermore, there appears to be a solid solution join from (Na,Li)<sub>7</sub>Th<sub>6</sub>F<sub>31</sub> to α-NaThF<sub>5</sub> or a new polymorph of Na<sub>7</sub>Th<sub>6</sub>F<sub>31</sub>, NLT being the phase limit of LiF solid solution. In fact, this latter possibility provides the most credible explanation for the existing phase equilibrium data in this region of the ternary diagram (Thoma, 1966). We propose that compositions on the low-LiF end of the solid solution join are achieved by simple replacement of lithium in both of its positions by sodium, resulting in a direct way in a composition limit Na<sub>7</sub>Th<sub>6</sub>F<sub>31</sub>. If, however, the composition limit is NaThF<sub>5</sub>, it is necessary to construct a more complicated mechanism, e.g. loss of Li(1) from its octahedral position and replacement of lithium by sodium in the general Li(2) position,

accompanied by a loss of fluorine, presumably from the polyhedron surrounding the ten-coordinate thorium, Th(2).

Figs. 1 through 8 are stereoscopic drawings (Johnson, 1965) of the contents of one asymmetrical unit and the coordination polyhedra around Th(1), Th(2), Na(1), Na(2), Na(3), Li(1), and Li(2) respectively. The drawing in Fig. 1 shows the unit cell rotated  $-30^{\circ}$  around  $a_1$  and all of the other Figures are rotated 45° around  $a_1$ . The choice of the F<sup>-</sup> nearest neighbors for Na(1), Na(3) and Li(2) is arbitrary and all of the fluorine ions less than 3.50 Å distant were included. It should be noted that the four nearest neighbor F<sup>-</sup> ions to Li(2) are at the corners of an irregular tetrahedron.

The coordination about Th(1) is the tripyramidal configuration normal for nine-coordinate atoms (Duffey, 1951; Kepert, 1965). The Th(2) environment may also be regarded as derived from the tripyramidal, by the insertion of an additional (fourth) fluorine ion in the central girdle.

The important interatomic distances not corrected for thermal motion are listed in Table 3. In the case of Li(2), corrections for independent thermal motion were



Fig.4. Coordination polyhedron around Na(1).

calculated; the results corresponded to Li(2)-F bond distances increased by 0.02-0.08 Å, but we are not confident that the large anisotropic thermal motion parameter for Li(2) calculated from the neutron data has any meaning. It is possible that it reflects a positional disorder.

For the most part, interatomic distances appear to be normal. Ligancy corrected Th-F distances are expected to be nearly 2.44 Å for 9- and 10-coordinated thorium (Pauling, 1960), and three thorium-fluorine distances are very significantly shorter than this: Th(1)-F(7), Th(1)-F(4), and Th(2)-F(8). However, the coordination polyhedra about both thorium atoms are quite irregular, and this should relax ligancy effects and packing constraints. The six short Na(1)-F(2) and Na(1)-F(10) distances correspond to an octahedral configuration about the sodium ion, for which an Na-F distance of 2.31 Å is expected (Pauling, 1960). There are no unusual F-F distances.

The observed X-ray structure factors and calculated structure factors from the XO parameters are listed in Table 4. The discrepancy index  $R = \Sigma ||F_o| - |F_c||/\Sigma/|F_o|$ is 0.0596 for all X-ray reflections and the standard deviation of an observation of unit weight,  $[\Sigma w(F_o - F_c)^2/(n_o - n_v)]^{1/2}$  is 0.80 ( $n_o$  is the number of reflections and

 $n_0$  the number of variables). The discrepancy index for XN (Table 2) is 0.0598.

Table 3. Interatomic distances for (Na,Li)<sub>7</sub>Th<sub>6</sub>F<sub>31</sub> determined form the XO and NO parameters (Table 2)

The number in parentheses is the standard error corresponding to the last digit.

	-					-		
	xo	NO		xo	NO		XO	NO
Tb(1)-F(7)	2.23(3)4	2.277(6)A	Th(2)-F(6)	2.25(4)A	2.425(6)A	L1(2)-F(8)	1.77(9)	1.92(3)
-2(9)	2.34(4)	2.351(6)	-7(8)	2.26(3)	2.293(6)	-7(6)	2.27(9)	2.03(3)
-F(4)	2.35(4)	2.297(6)	-7(3)	2.32(4)	2.378(6)	-7(1)	2.33(9)	2.08(4)
-\$(3)	2.38(4)	2.371(6)	-7(4)	2.36(4)	2.383(6)	-7(5)	2.47(9)	2.75(3)
-7(5)	2.38(4)	2.352(5)	-7(1)	2.40(4)	2.386(6)	-7(3)	2.49(9)	2.77(4)
-1(2)	2.41(4)	2.380(6)	-F(10)	2.41(4)	2.355(6)	-1(9)	2.62(9)	3,18(3)
-7(1)	2.45(4)	2.522(6)	-7(9)	2.43(3)	2.510(6)	-1(7)	2,70(9)	1.98(3)
-7(10)	2.49(4)	2.435(6)	-2(6)	2.44(4)	2.573(6)	-7(4)	2.89(9)	3.37(3)
-7(5)	2.63(4)	2.354(6)	-7(2)	2.46(4)	2.425(6)			
			-7(11)	2.63(2)	2.514(3)			
N=(1)-3F(10)	2.26(5)	2.26(2)	Ma(2)-3P(9)	2.25(4)	2.30(2)	Na(3)-3P(7)	2.56(6)	3.00(2)
-37(2)	2.27(4)	2.29(1)	-3F(1)	2.30(4)	2.299(8)	-37(4)	2.69(4)	2.730(6)
-3F(6)	3.04(6)	3,41(2)	-37(5)	3.34(6)	3.07(2)	-32(3)	2.78(4)	2.649(6)
-7(11)	3.1(2)	3.18(3)	L1(1)-3P(8)	2.0(1)	2.06(3)	-37(8)	2.96(6)	2.50(2)
-7(11)	3.5(2)	3.46(3)	-37(7)	2.1(1)	1.96(2)			
F(1)-F(2)	2.53(4)	2.517(5)	¥(1)-¥(9)	2.64(6)	2.717(8)	P(1)-P(5)	2.72(6)	2.838(8)
F(1)-F(7)	2.72(6)	2.688(7)	¥(1)-F(5)	2.62(6)	2.593(7)	F(1)-F(3)	2.83(6)	2.763(8)
F(2)-F(4)	2.90(5)	2.915(8)	P(2)-P(7)	2.88(5)	3.141(8)	F(2)-F(8)	2.76(5)	2.817(8)
P(2)-F(10)	2.66(6)	2.623(8)	F(2)-F(6)	2.51(5)	2.696(8)	F(2)-F(6)	2.80(5)	2.761(8)
F(3)-F(4)	2.65(6)	2.718(9)	F(3)-F(7)	2.72(6)	2.702(8)	7(3)-7(5)	2.96(6)	2.893(8)
F(3)-F(9)	2.67(5)	2.713(8)	F(3)-F(4)	2.86(7)	2.724(9)	P(3)-P(9)	2.67(6)	2.724(8)
F(3)-F(8)	2.94(5)	2.714(8)	P(4)-P(7)	2.74(5)	2.942(8)	F(4)-F(10)	2.77(5)	2.723(8)
F(4)-F(10)	2.66(6)	2.712(8)	F(4)-F(8)	2.59(6)	2.669(8)	7(4)-7(6)	2.76(6)	2.812(8)
F(5)-F(10)	2.64(6)	2.794(8)	F(5)-F(7)	2.68(5)	2.828(7)	7(5)-7(9)	3.05(5)	2.726(8)
F(6)-F(11)	1.82(7)	2.416(8)	F(6)-F(10)	2.78(6)	3.041(8)	F(6)-F(6)	2.71(7)	3.88(2)
F(6)-F(6)	2.71(7)	3.88(2)	F(6)-F(9)	2.89(6)	2.593(9)	F(6)-F(8)	2.76(5)	2.638(7)
F(7)-F(8)	2.80(4)	2.757(6)	P(7)-F(7)	2.94(5)	2.921(9)	P(7)-F(7)	2.94(5)	2.921(9)
F(7)-F(8)	2.91(4)	2.849(6)	F(8)-F(8)	2.93(5)	2.816(9)	7(8)-7(8)	2.93(5)	2.816(9)
F(9)-F(10)	2.53(4)	2.503(5)	F(10)-F(11)	2.26(5)	2.26(2)	F(10)-F(10)	3.00(7)	3.24(2)
F(10)-F(10)	3,00(7)	3,24(2)						



Fig. 5. Coordination polyhedron around Na(2).



Fig. 6. Coordination polyhedron around Na(3).

The final neutron structure factors appear in Table 5, in which  $|FOBS|^2 = 10^{24} |F_o|^2$  cm<sup>2</sup>. These data correspond to the NO parameters listed in Table 2 and yield for all data:

and

$$R_{1} = \Sigma ||F_{o}| - |F_{c}|| / \Sigma |F_{o}| = 0.096$$
$$R_{2} = \Sigma ||F_{c}|^{2} - |F_{c}|^{2} |\Sigma |F_{c}|^{2} = 0.110$$

# Table 4. The observed and calculated structure factors and $\alpha$ angles for X-ray data

 $\alpha \times 10$ 



Actual least-squares adjustments, however, minimized the function  $\Sigma w ||F_o|^2 - |F_c|^2|^2$ , in which the weights, w, were taken as  $1/\sigma^2$  (F<sub>o</sub>) if  $|F_o|^2 \ge 2\sigma$ , and zero if  $|F_o|^2 < 2\sigma$ . For the data greater than  $2\sigma$ , the R values defined above become:  $R_1 = 0.061$ ,  $R_2 = 0.091$ . The standard deviation of an observation of unit weight,

$$\sigma = [\Sigma w (|F_o|^2 - |F_c|^2)^2 / (n_{\text{obs}} - n_{\text{var}})]^{1/2}$$

is 1.25 for the data greater then  $2\sigma$ , and 1.17 for all data. The discrepancy indices  $R_1$  and  $R_2$  for NX (Table 2) are respectively 0.1096 and 0.1279.

We owe a great debt to C.K. Johnson for supervising the data collection on the Oak Ridge Automatic Three Circle Neutron Diffractometer, and for introducing one of us (DRS) to neutron diffraction methods. We



Fig.7. Coordination polyhedron around Li(1).

thank H.A. Friedman for supplying the lithium-7 material. J.C.Ogle and H.G. Smith graciously provided X-ray and neutron Laue equipment. L.E. Alexander made available to us the Mellon Institute piezoelectric detector. We appreciate the time D.P. Shoemaker spent in reviewing the original draft.

# Table 5. The square of the observed and calculated structure factors and the $\alpha$ angles for neutron data $|FOBS|^2 = 10^{23} |F_o|^2 (cm)^2$ .

The tabulated values have been multiplied by 10.

and a state a strate a strate and a strategies an \$2956219389122823643555 第8年15月1日、191 193: °26X:555\* 1759 206 1656 75 14 1574 883 1461 153 1618 227 87 ៶៰៝៝៝៝៵៶ៜ៵៹ៜ៵៹៹៵៶៵៹៵៶៵៹៶៵៹៵៰ឨ៵ឨ៵ឨ៵ឨ៵៹ 125 51 1139 1779 1655 31 273 \*282333282888822<sup>0</sup> \*\$2527292331<sup>0</sup>13228673 11 12 13 14 15 16 17 8 10 12 14 15 18 20 1760 126 205 156 1513 156 1613 130 177 723 1642 1692 923 1764 1613 310013777820557672133412537821613100318¥18546500476711160018257161217721030526037225717574271257427124566616586107821351516205139410 0-234567890123115678 <sup>°</sup> <u>ㅋ</u>Ĕ┇¥¥65568<sup>°</sup> 의사82×\*\*9° 8138¥\$618138495<sup>°°</sup> 22¥¥8662254<sup>°</sup> 7288193777<sup>°°</sup> 7598\*78¥656677862788 10 12 14 16 137529292923373 1713 1558 232 1771 1051 1704 1624 983 7.645423261 02366303333272 2969336799 22281557655 141 131 1518 1593 1274 1411 639 13 14 15 16 17 18 19 521 749 1314 902 1626 351 77 1071 921 606 957 16 1455 1732 166 1322 828×8×82 204 721 721 721 166 180 02.20 59 20 57 961 29 T 461 71 453 102 1



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Fig.8. Coordination polyhedron around Li(2).

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# The Crystal and Molecular Structure of Vanadyldeoxophylloerythroetioporphyrin-1,2-dichloroethane Solvate: an Analogue of Chlorophyll\*

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Single crystals of vanadyl deoxophylloerythroetioporphyrin-1,2-dichloroethane solvate,  $C_{32}H_{34}N_4VO.C_2H_4Cl_2$ , were grown by a vapor-phase diffusion technique. Pertinent crystal data are: space group  $P_{21}/a$ , a=17.963 (3), b=13.961 (1), c=14.090 (3) Å,  $\beta=116.97$  (1)°, Z=4. The structure was solved by locating the position of the vanadium atom from a three-dimensional Patterson map and calculating successive Fourier series to locate the remaining atoms. A difference map revealed positions for 28 of the 38 hydrogen atoms, and these were used in the refinement procedure. Full-matrix leastsquares calculations gave a final weighted R of 5.5% for approximately 2000 observed and unobserved reflections. The interesting feature of this molecule is the effect of the fifth isocyclic ring, E, on the porphyrin skeleton. Results indicate that this ring introduces strain into the molecule and causes distortions in the bond angles and most likely causes a shortening of the V–N bond to the ring C atom from an average of 2.10 Å for the other three V–N bonds to 1.96 Å for the shortened bond. The carbon and nitrogen atoms of the porphyrin skeleton are relatively planar and the vanadium atom lies 0.48 Å out of the plane of the four nitrogen atoms.

### Introduction

In recent years new insights have been gained into the structure of chlorophyll and its related compounds. The basic structure of chlorophyll was worked out in the late 1930's by Hans Fischer (Fischer & Wenderoth, 1940) and subsequently confirmed through a very elaborate synthesis by R.B. Woodward and associates (Woodward, 1961). The accumulation of chemical information has recently made possible the description of the absolute configuration of chlorophyll a (Fleming, 1967). Crystallographers have shown an interest in an accurate structure determination for chlorophyll a, but the obstacle has so far been one of obtaining suitable crystals. At Mellon Institute an investigation has been begun into the feasibility of growing crystals of mater-

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